

Tok elektrického prúdu v nanoštruktúrach

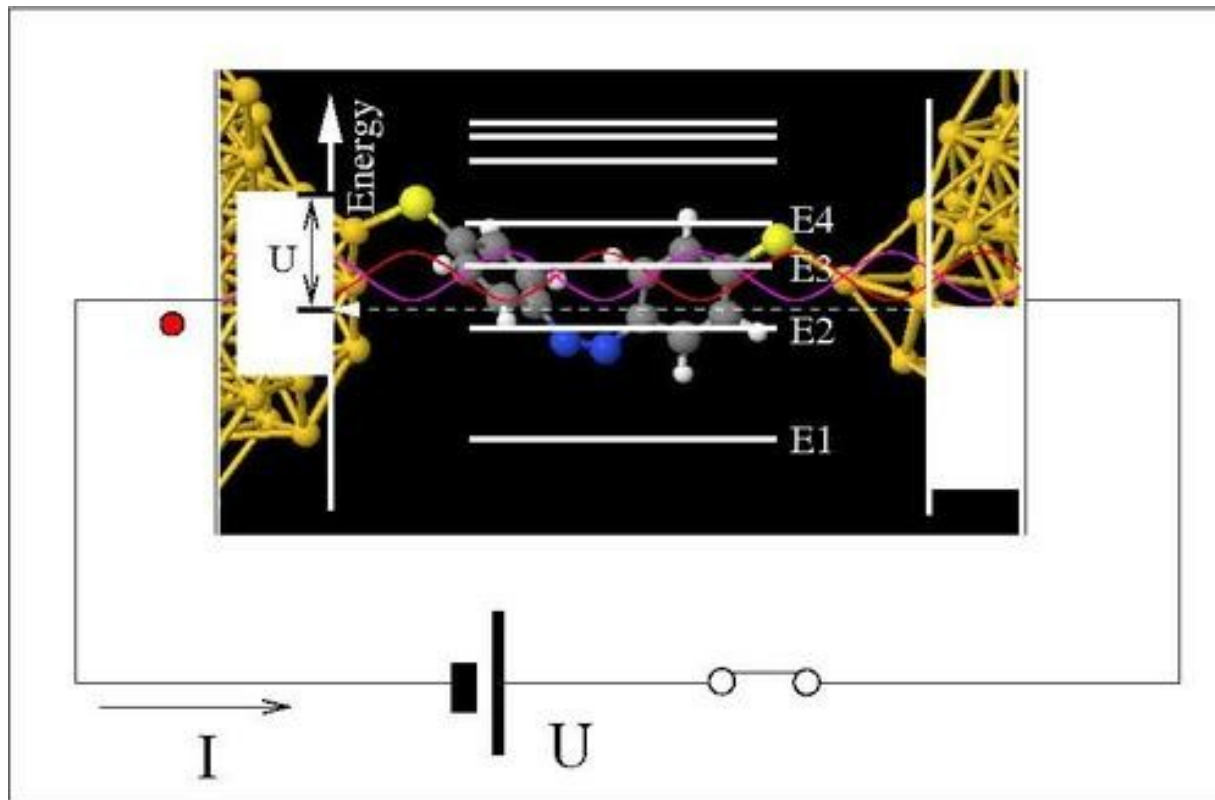
Habilitačná prednáška

Peter Bokes

Katedra fyziky, FEI STU



Tok elektrického prúdu v nanoštruktúrach



Ab-initio model

Nerovnovážny
kvantový proces

Kvalitatívne
porozumenie

Merateľné vlastnosti

$$I(U), \quad G = \left. \frac{dI}{dU} \right|_{U=0}$$

[Konduktancia]

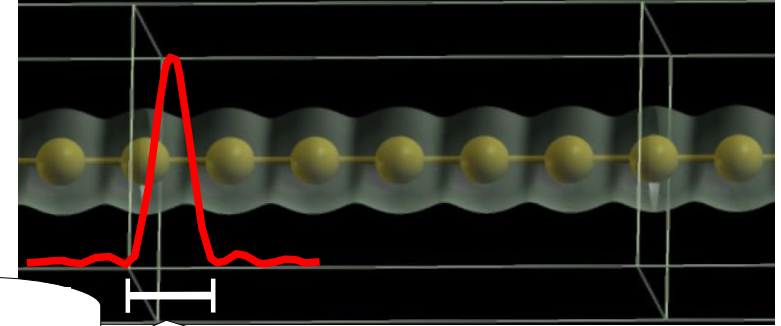
Osnova prednášky

- Pohyb elektrónu a vlnová funkcia – vlnový balík
- Pravdepodobnosť prechodu cez nanoštruktúru - transmisia
- Pohyb mnohých elektrónov – ortogonálne vlnové balíky
- Landauerov vzťah pre elektrický prúd cez nanoštruktúru
- Transport cez realistické nanoštruktúry

Vlnový balík

- Časový vývoj stavu – Schrodingerova rovnica

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = \hat{H} \psi(x, t)$$



- Stacionárne stavy:

$$\hat{H} \phi_k(\vec{r}) = E_k \phi_k(\vec{r})$$

$$\phi_k(x) = A e^{ikx}, \quad E_k = \frac{\hbar^2 k^2}{2m}, \quad p = \hbar k$$

$$\hat{H} = -\frac{\hbar^2}{2m_e} \frac{d^2}{dx^2} + u_e(x)$$

$\Delta x \Delta k \geq 4\pi$ “Heisenbergov princíp neurčitosti”



$$\psi_k(x, t) = \phi_k(x) e^{-iE_k t / \hbar}$$

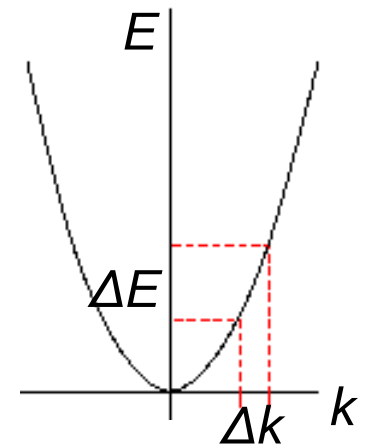
$$P(x) = |\psi_k(x, t)|^2 = |A|^2$$

- Vlnový balík:** súčet (integrál) vlastných stavov

$$\psi_{E, \Delta E}^{R/L}(x, t) = \int_{k-\Delta k/2}^{k+\Delta k/2} dk A e^{i(\pm kx - E_k t / \hbar)}$$

- lokalizovaný v priestore, a pohybuje sa **grupovou rýchlosťou**

$$P(x) = |\psi_{E, \Delta E}^R(x, t)|^2 = 4|A|^2 \frac{\sin^2\left(\frac{\Delta k}{2}(x - vt)\right)}{(x - vt)^2}, \quad v = \frac{1}{\hbar} \frac{dE_k}{dk}$$

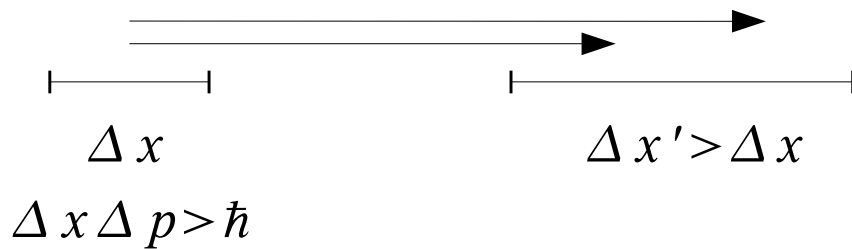


- normovateľný na 1** (pre každé $t > 0$) ak

$$\int_{-\infty}^{\infty} dx P(x) = 1 \Rightarrow A = (2\pi \Delta k)^{-1/2}$$

Disperzia vlnového balíka

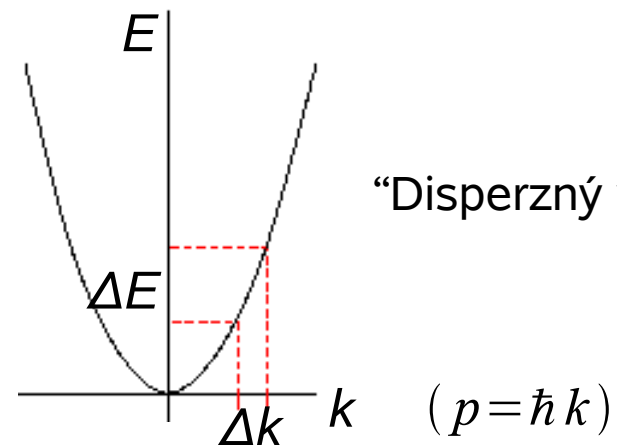
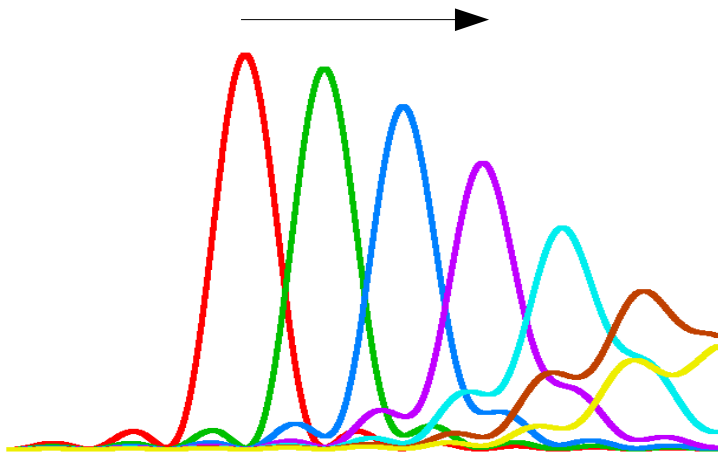
- **Disperzia** - nevyhnutná v dôsledku Heisenbergovho princípu neurčitosti:



$$\Delta x' = \Delta x + \Delta v t, \Delta v = \frac{\hbar}{\Delta x m}$$

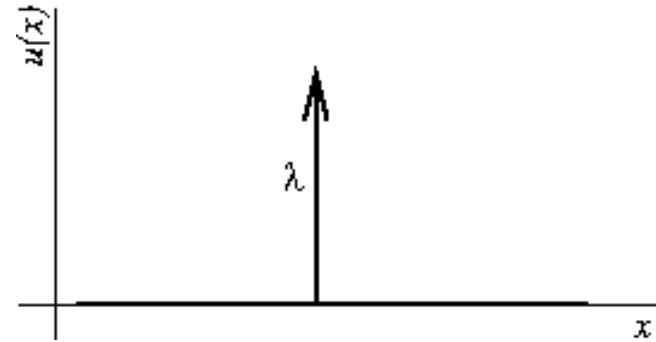
$$v = \frac{\partial E(k)}{\hbar \partial k}, \Delta v = \frac{\partial^2 E(k)}{\hbar \partial^2 k} \Delta k = \frac{\Delta p}{m}$$

keď $E(k)$ nie je lineárna funkcia
alebo pre $\Delta k/k \gg 1$



Transmisný koeficient, transmisia

- Hamiltonián: $\hat{H} = -\frac{\hbar^2}{2m_e} \frac{d^2}{dx^2} + u_{eff}^\lambda(x)$



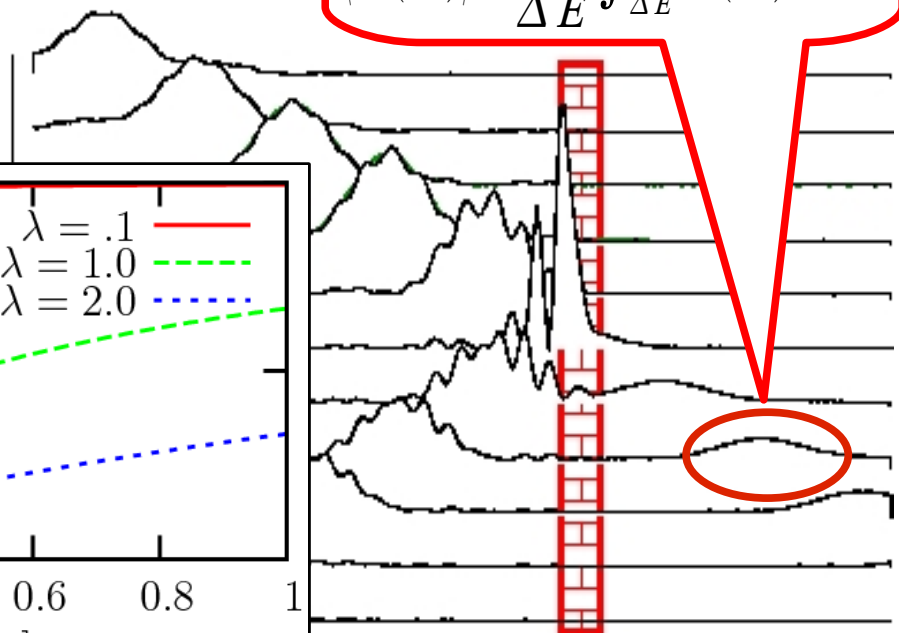
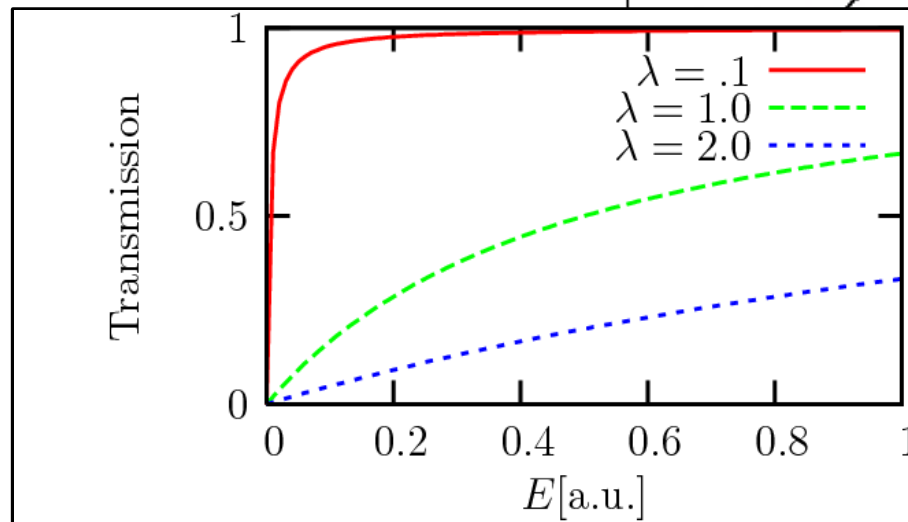
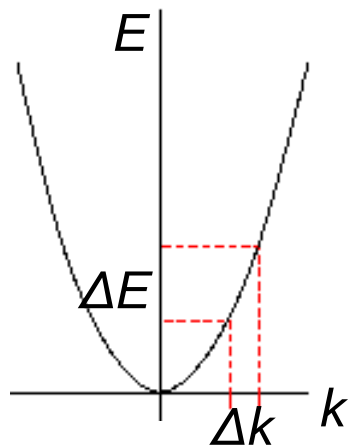
- Stacionárne riešenia pre $E_k = \hbar^2 k^2 / (2m_e)$

$$\phi_k^R(x) = A \begin{cases} e^{ikx} + r_k e^{-ikx} & x \ll 0 \\ t_k e^{ikx} & x \gg 0 \end{cases} \quad \phi_k^L(x) = A \begin{cases} \tilde{t}_k e^{-ikx} & x \ll 0 \\ \tilde{r}_k e^{-ikx} & x \gg 0 \end{cases}$$

Transmisia: $T(E) = |t_k|^2$
 $\langle T(E) \rangle = \frac{1}{\Delta E} \int_{\Delta E} T(E) dE$

- Balík

$$\psi_{E, \Delta E}^R(x, t) = \int_{k-\Delta k/2}^{k+\Delta k/2} A \phi_k^R(x) e^{-iE_k t / \hbar} dk$$



Transmisný koeficient, rezonancia

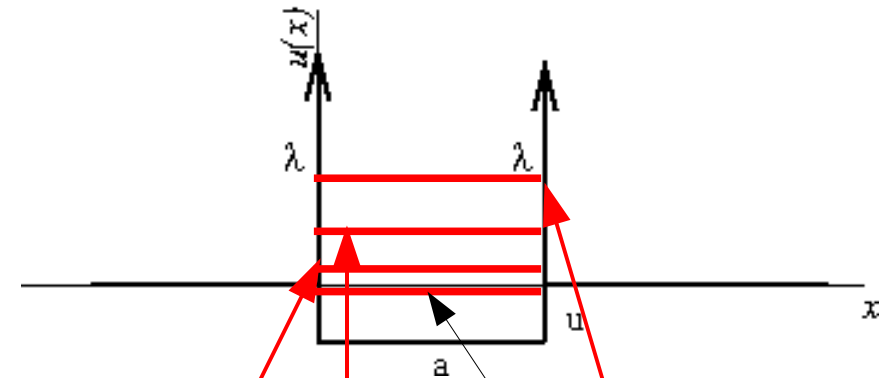
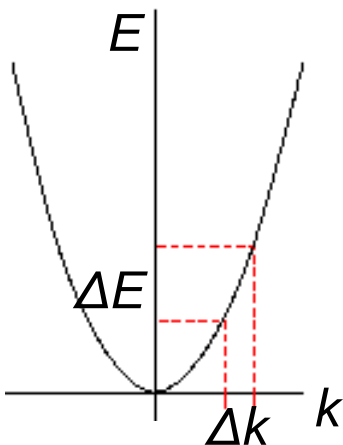
- Hamiltonián: $\hat{H} = -\frac{\hbar^2}{2m_e} \frac{d^2}{dx^2} + u_{eff}^\lambda(x)$

- Stacionárne riešenia pre $E_k = \hbar^2 k^2 / (2m_e)$

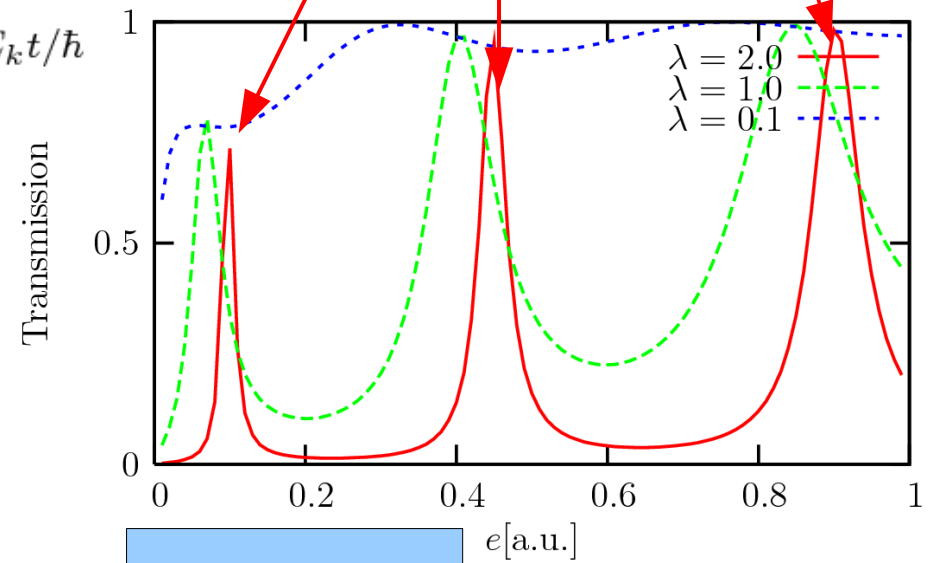
$$\phi_k^R(x) = A \begin{cases} e^{ikx} + r_k e^{-ikx} & x \ll 0 \\ t_k e^{ikx} & x \gg 0 \end{cases}$$

- Balík

$$\psi_{E,\Delta E}^R(x,t) = \int_{k-\Delta k/2}^{k+\Delta k/2} A \phi_k^R(x) e^{-iE_k t/\hbar}$$



Viazaný stav



Rezonancia

Kvantový popis mnoho elektrónov

- **Nestacionárna teória funkcionálu hustoty** [Time-dependent density-functional theory, TDDFT, E. Runge, E. K. U. Gross (1984)] – elektróny ako fiktívne neinteragujúce častice, pohybujúce sa v efektívnom potenciáli.

- Elektróny sa nachádzajú v stavoch daných **ortogonálnymi** vlnovými funkciami $\psi_n(\vec{r}, t)$

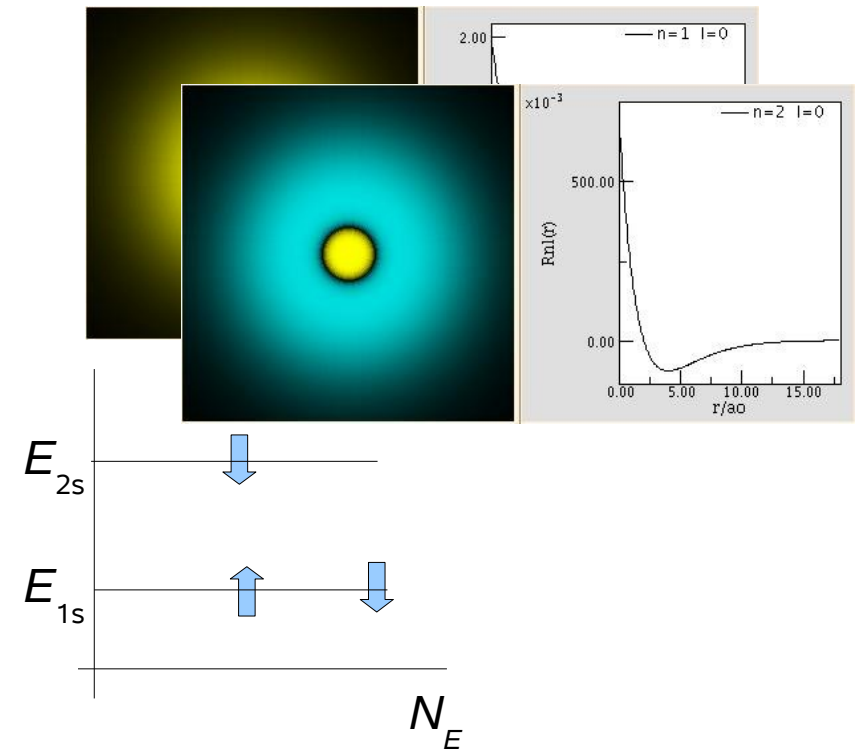
$$\int_{-\infty}^{+\infty} dx \psi_n^*(x) \psi_m(x) = \delta_{nm}$$

- **Pauliho princíp** – najviac 2 elektróny v jednom stave, potom *exaktná* hustota bude

$$n(\vec{r}, t) = \sum_n N_n |\psi_n(\vec{r}, t)|^2, \quad N_n \in \{0, 1, 2\}$$

- Vlnové funkcie sa menia s časom podľa rovnice

$$i\hbar \frac{\partial}{\partial t} \psi_n(\vec{r}, t) = \hat{H} \psi_n(\vec{r}, t), \quad \hat{H} = -\frac{\hbar^2}{2m_e} \Delta + u_{eff}[n(\vec{r}', t')](\vec{r}, t)$$



1D kvantový vodič

- stacionárna Schrödingerova rovnica

$$\hat{H}\phi_n(\vec{r}) = E_n\phi_n(\vec{r}), \quad \hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

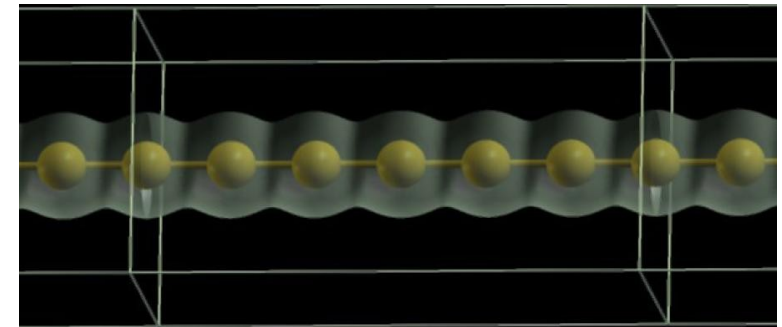
$$\phi_n(x) = \frac{1}{\sqrt{L}} e^{ik_n x}, \quad E_n = \frac{\hbar^2 k_n^2}{2m}$$

$$k_n = \frac{2\pi}{L} n, \quad n = 0, \pm 1, \pm 2, \dots$$

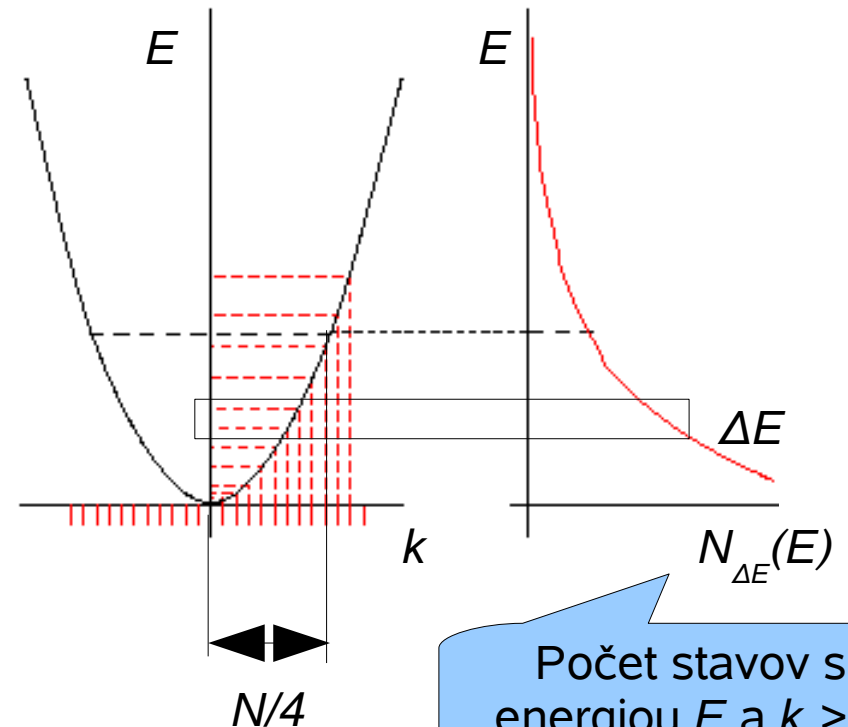
- hustota, energia

$$n(x) = \sum_{|n| < N/4} 2|\phi_n(x)|^2 = N/L$$

$$E = \sum_{|n| < N/4} 2E_n = \dots$$



L, N elektrónov



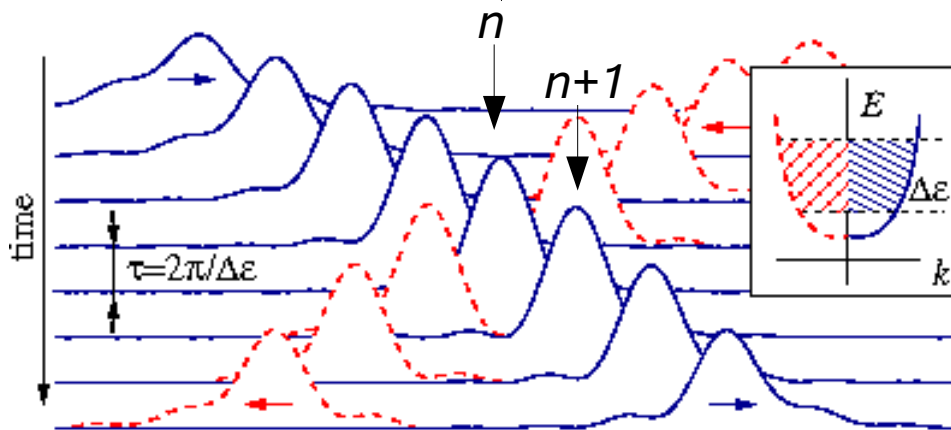
Počet stavov s energiou E a $k_n > 0$

Ortogonalne vlnové balíky

• Balíky

$$\psi_{E,\Delta E,n}^{R/L}(x) = \frac{\hbar}{\sqrt{2\pi m_e \Delta E}} \int_{k-\Delta k/2}^{k+\Delta k/2} dk \sqrt{k} \phi_k^{R/L}(x) e^{-iE_k t/\hbar} \Big|_{t=n\tau}$$

“Stroboskopické balíky”



s časovým krokom $\tau = \frac{2\pi\hbar}{\Delta E}$

sú ortogonálne:

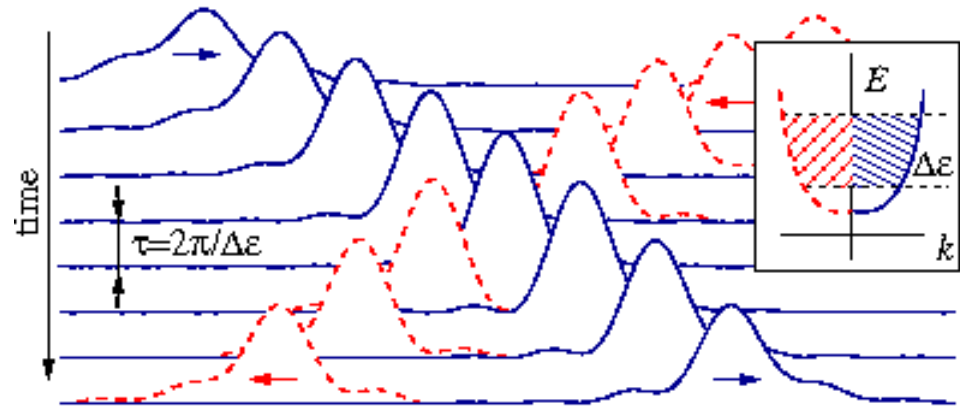
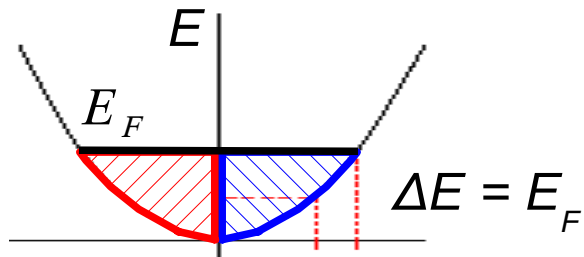
$$\int_{-\infty}^{+\infty} dx \psi_{E,\Delta E,n}^*(x) \psi_{E,\Delta E,m}(x) = \delta_{nm}$$

- “Princíp neurčitosti”: $\Delta E \tau \sim 2\pi\hbar$
- Navzájom ortogonálne sú aj balíky
 - s disjunktných intervalov energií
 - naopak idúce [“R” s “L”]
- Všetky balíky dokopy tvoria **úplný systém funkcií**

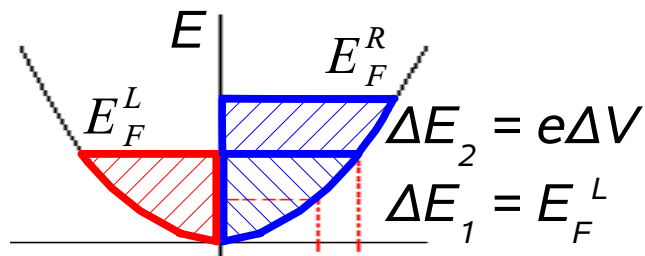
Obsadenie balíkov

- **Fermiho energia** - rezervoár elektrónov naplní všetky stavy s $E < E_F$

$$n(\vec{r}, t) = \sum_{occ} 2|\psi_n(\vec{r}, t)|^2$$

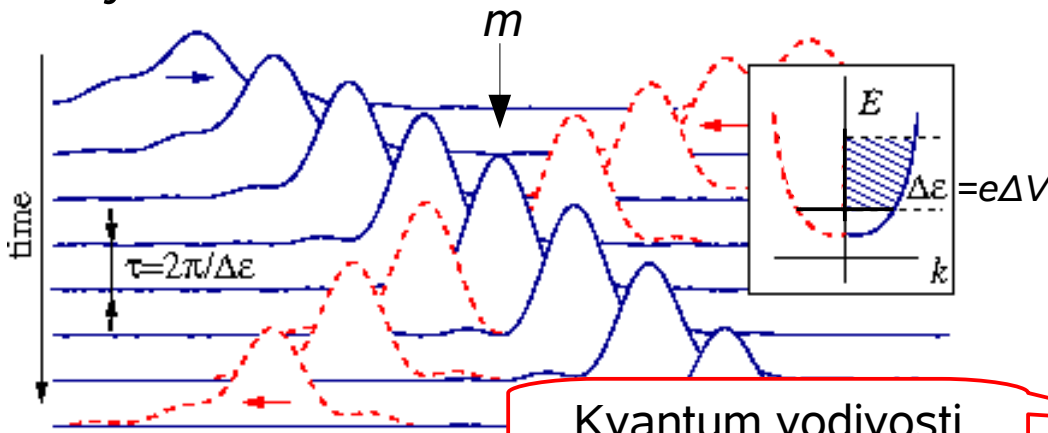


- **Aplikované napätie** – rozdiel Fermiho energií rezervoárov pripojených na vodič zľava a zprava $\Delta V = (E_F^R - E_F^L)/e$



Landauerov vzťah pre elektrický prúd

Čistý vodič:



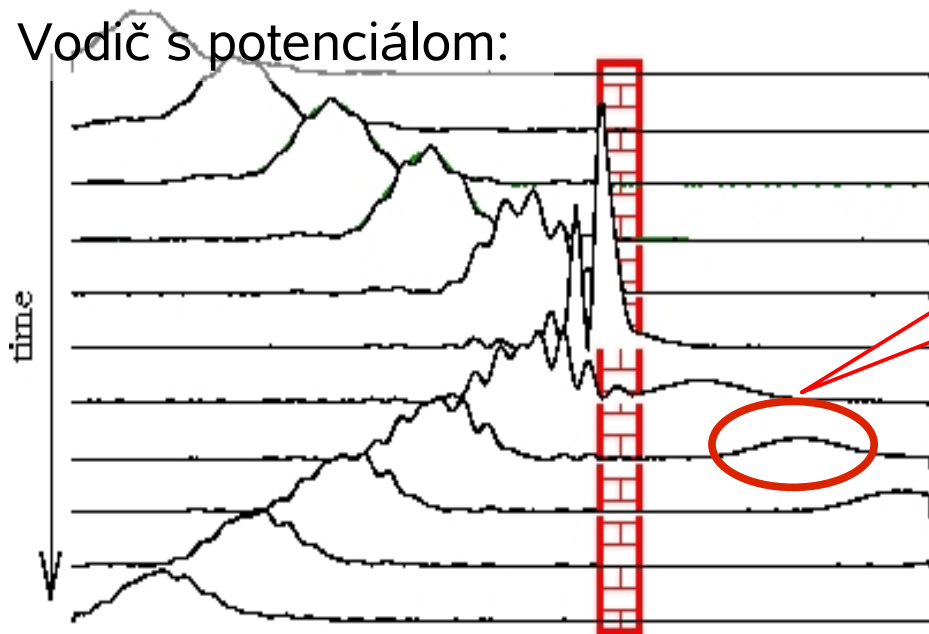
$$I_m(t) = \frac{e}{\tau} N_m(t)$$

$$I = \frac{e}{2\pi\hbar/(eV)} N_m(t) = \frac{2e^2}{h} V$$

$$G_0 = \frac{2e^2}{h}, R_0 = \frac{h}{2e^2} \approx 12.9\text{k}\Omega$$

Kvantum vodivosti

Vodič s potenciálom:



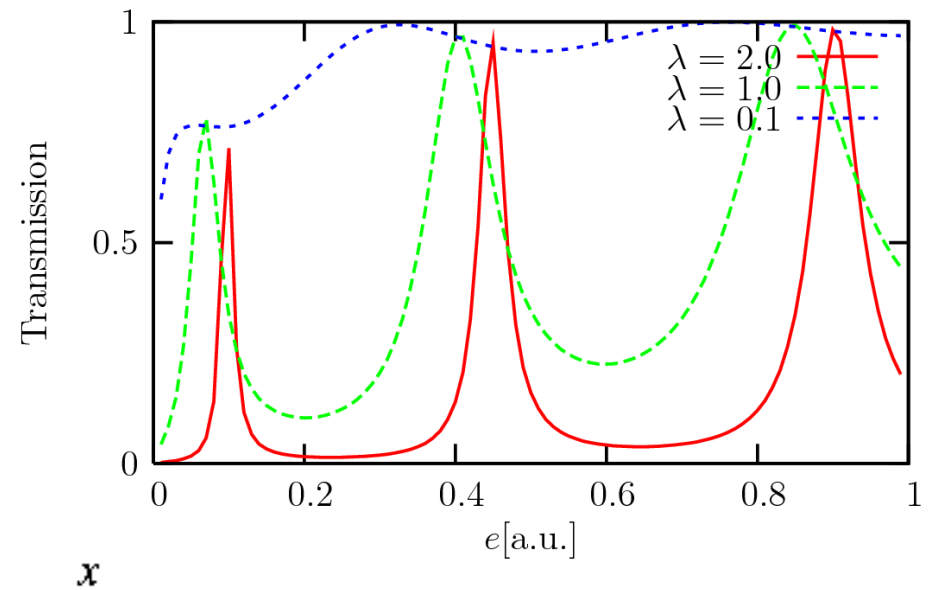
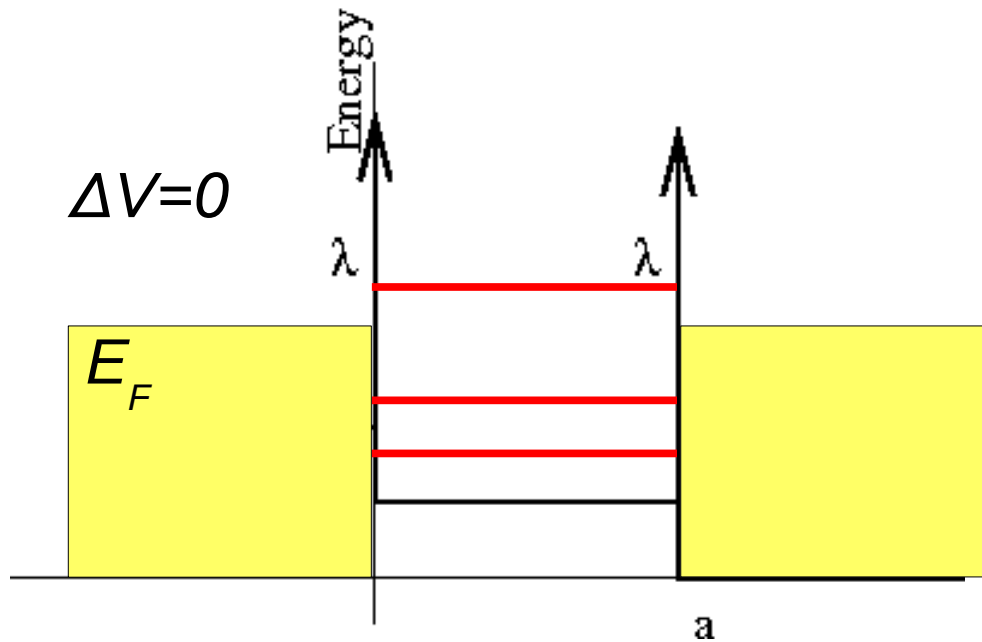
Transmisia dá obsadenie:
 $N_m(t) = 2\langle T(E) \rangle$

Viac kanálov

=> Landauerov vzťah

$$I = \frac{2e}{h} \int_{e\Delta V} \sum_i T_i(E) dE$$

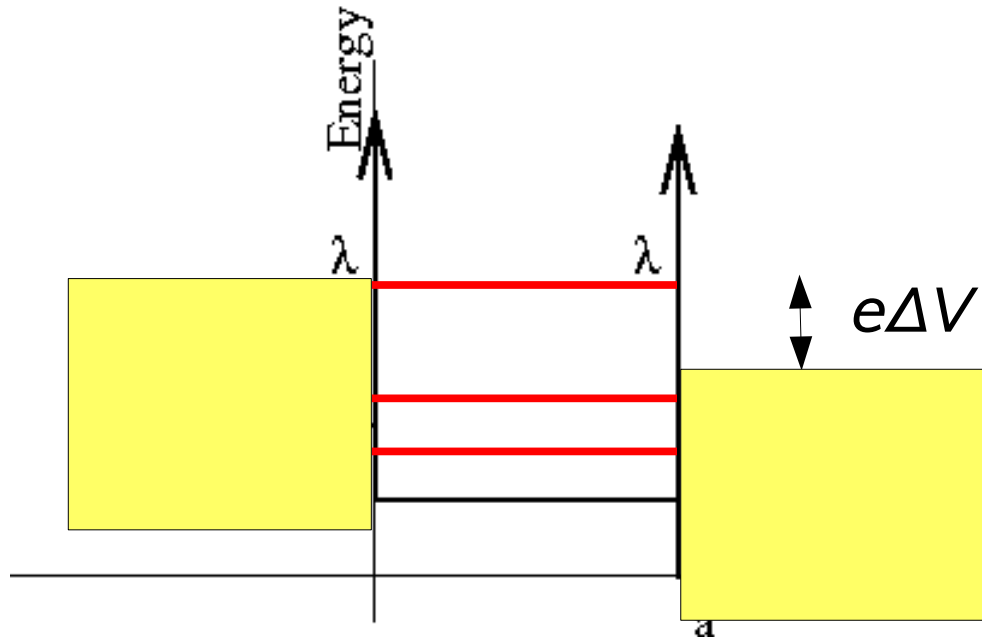
Rezonančné tunelovanie, I/V



Stacionárny prúd:

$$I = \frac{2e}{h} \int_{e\Delta V} T(E) dE$$

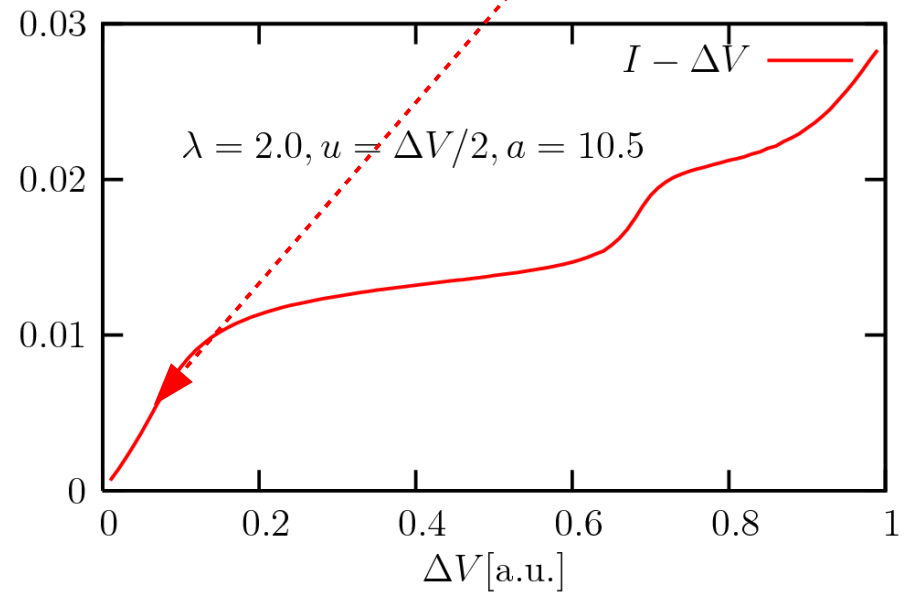
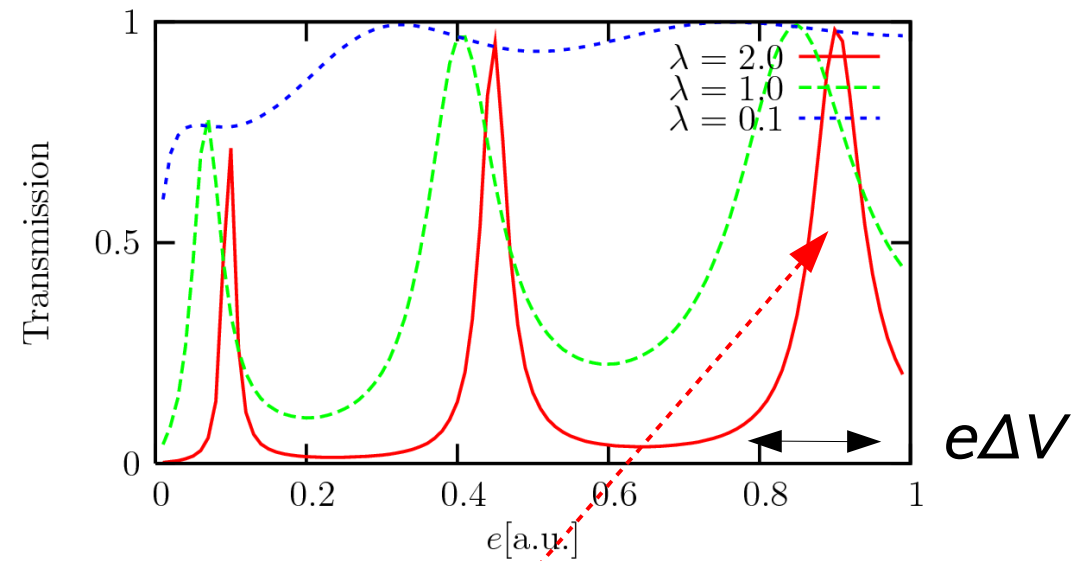
Rezonančné tunelovanie, I/V



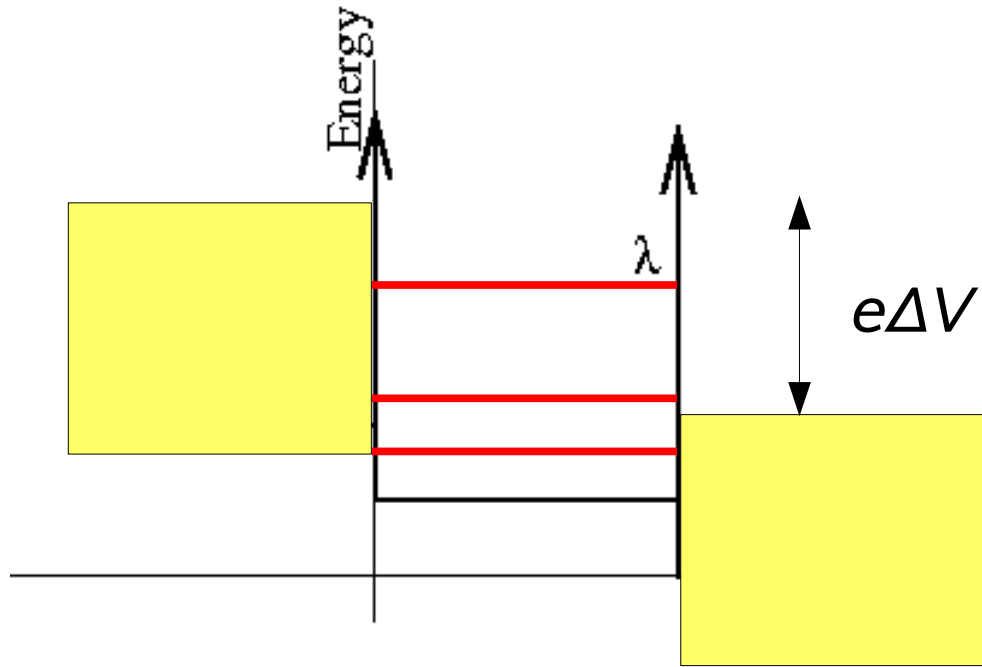
Stacionárny prúd:

$$I = \frac{2e}{h} \int_{e\Delta V} T(E) dE$$

I [a.u.]



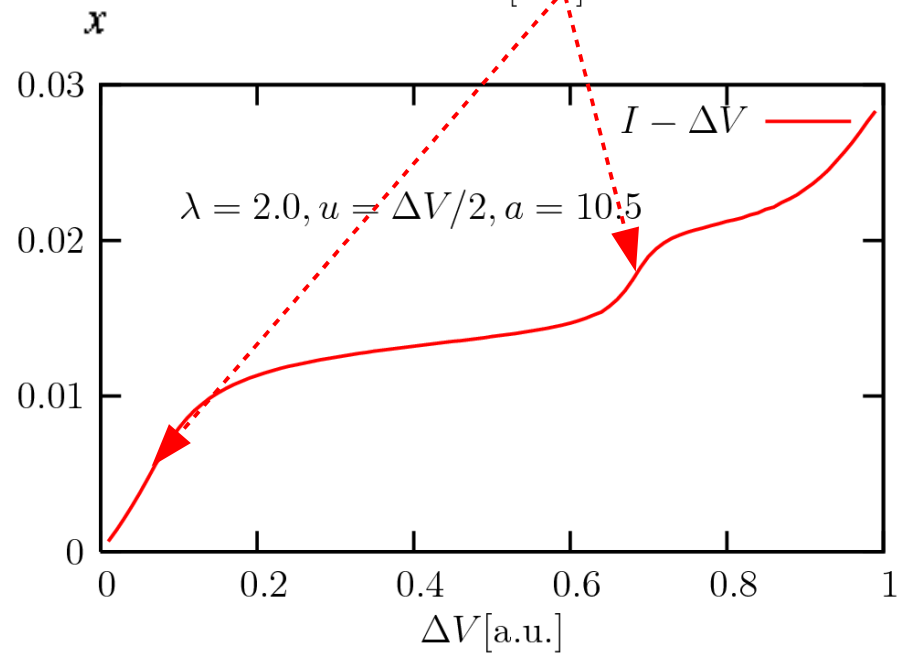
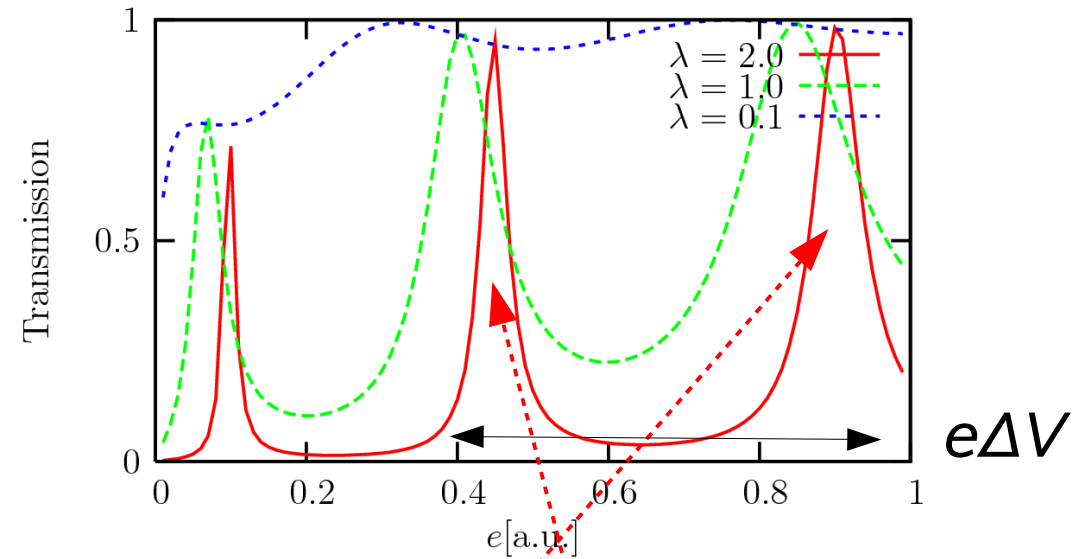
Rezonančné tunelovanie, I/V



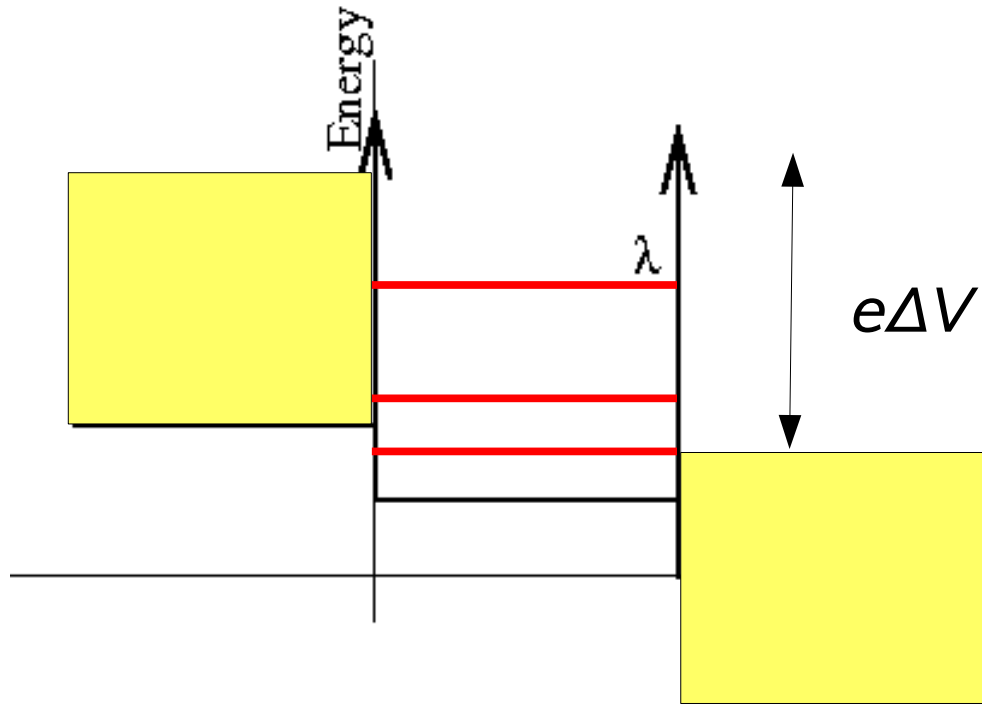
Stacionárny prúd:

$$I = \frac{2e}{h} \int_{e\Delta V} T(E) dE$$

I [a.u.]



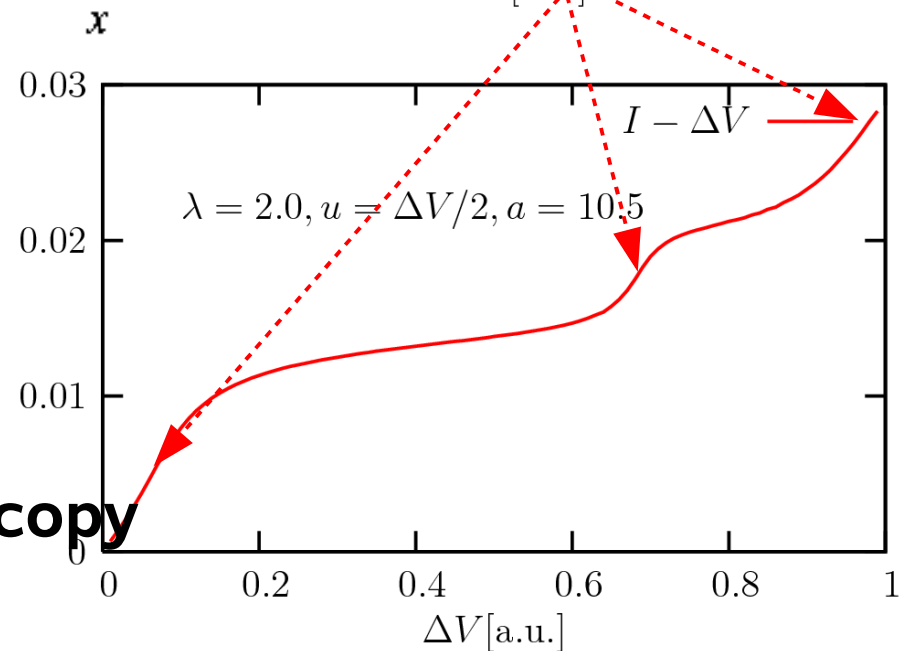
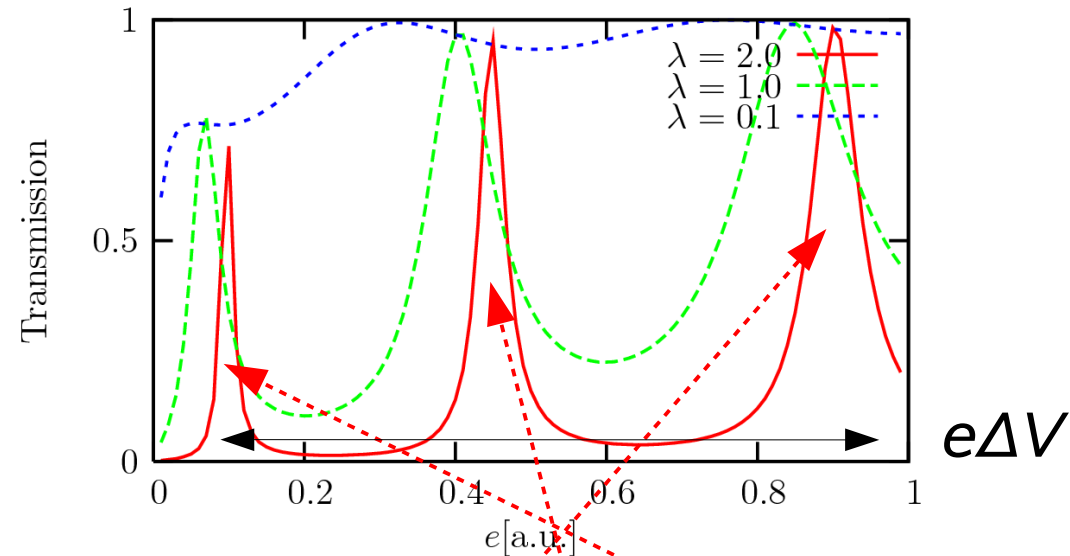
Rezonančné tunelovanie, I/V



Stacionárny prúd:

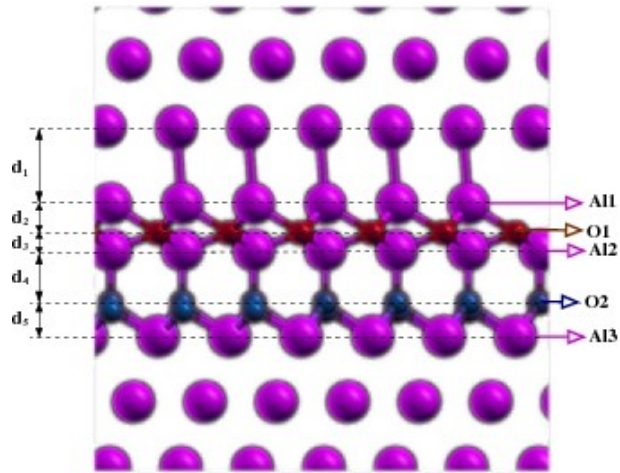
$$I = \frac{2e}{h} \int_{e\Delta V} T(E) dE \quad I[\text{a.u.}]$$

=> Scanning Tunneling Spectroscopy (STS/STM)

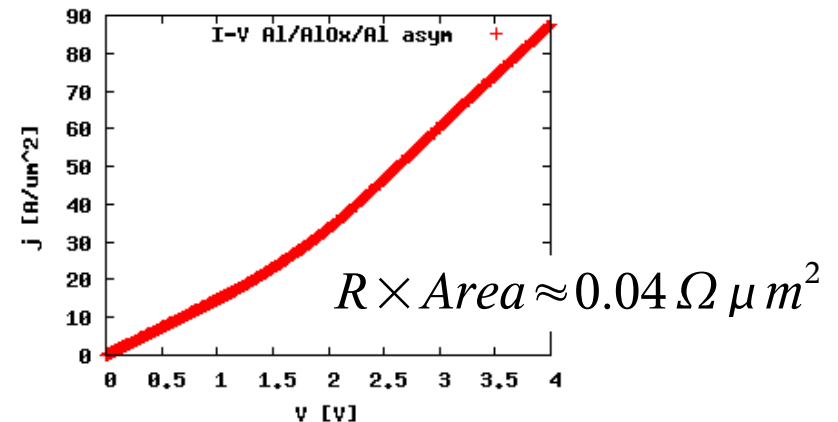
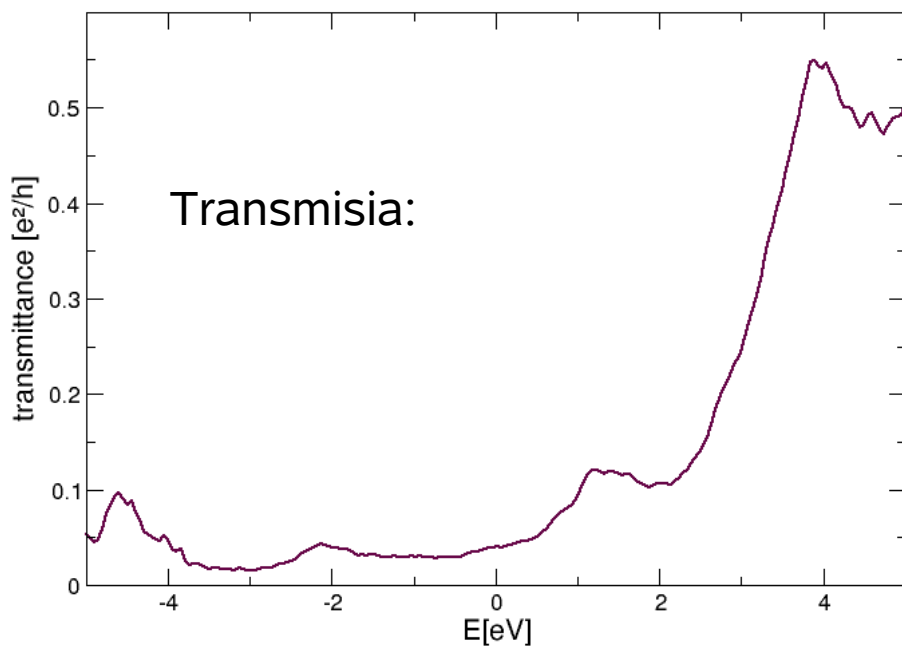
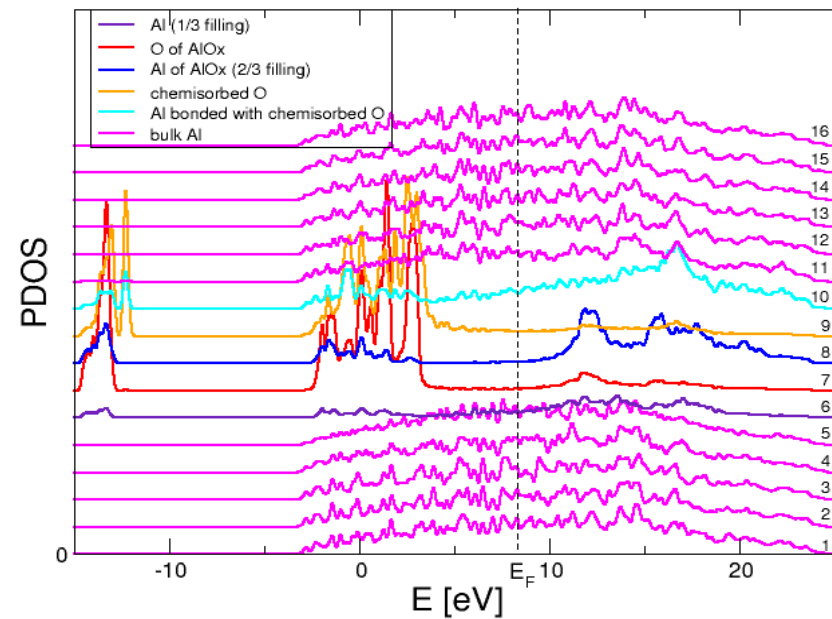


Transport cez rozhranie $Al/AlO_x/Al$

Atomárna geometria:

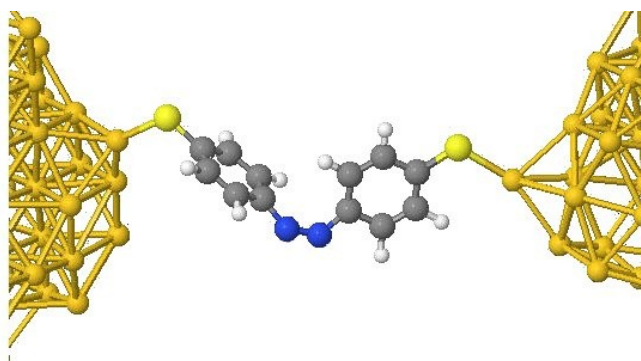


Lokálna hustota stavov:

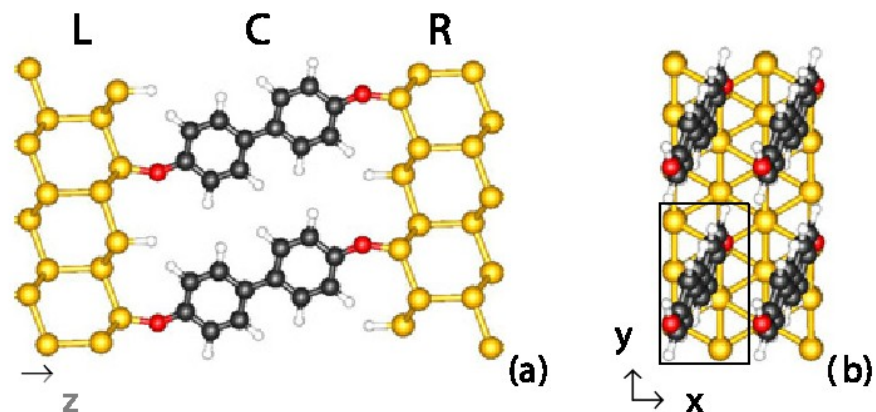


M. Diešková, A. Ferretti, PB (2009),
M. Diešková, M. Konôpka, PB (2007)

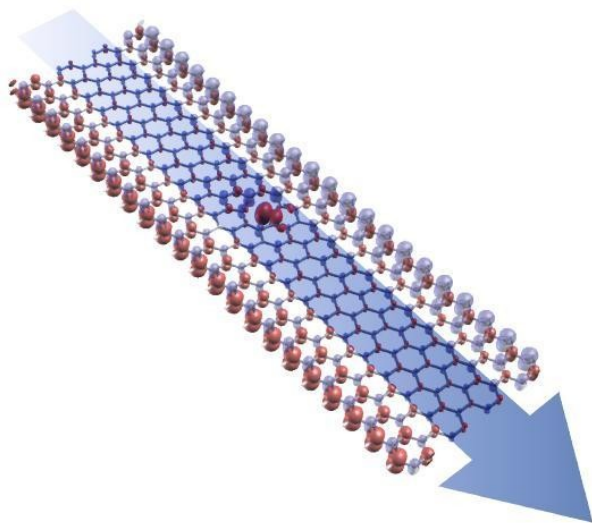
...a iné reálne nanokontakty



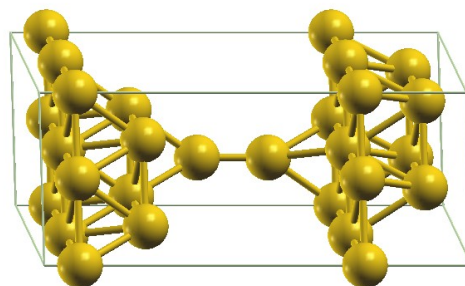
Au-Azobenzene-Au
Konôpka, Turanský, Štich (2008)



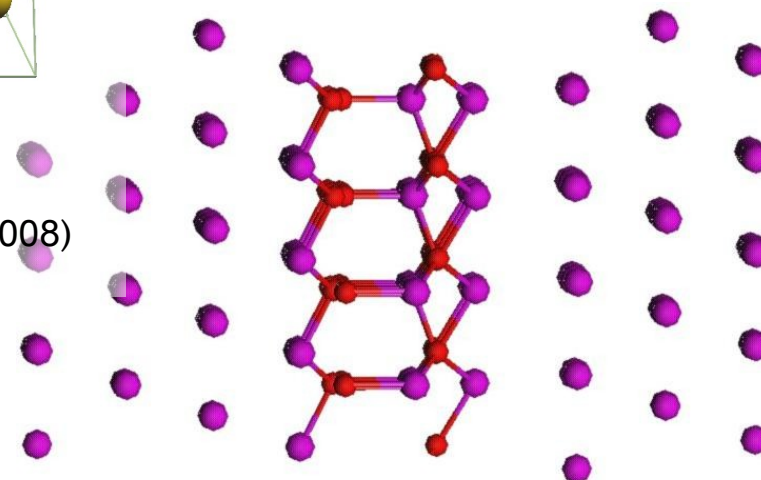
Si-C₁₂H₁₀O₂-Si
Bonferroni, Ferretti, Molinari (2008).



Graphene
Rignanese, Charlier (2008)



Au kontakt,
Verstraete, PB, Godby (2008)



Al-AlO_x-Al rozhranie
Diešková, Konôpka, PB (2007)

Zhrnutie

- Vlnový balík – vlnová funkcia dostáva zmysel
- Ortogonálne vlnové balíky pre mnoho elektrónov
- Landauerov vzorec pre elektrický prúd
- Reálne nanoštruktúry